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## Chemical and *in silico* tuning of the magnetisation reversal barrier in pentagonal bipyramidal Dy(III) single-ion magnets†

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**Two new air-stable axial Dy(III) single-ion magnets, [Dy(H<sub>2</sub>O)<sub>5</sub>-(HMPA)<sub>2</sub>]Cl<sub>3</sub>·HMPA·H<sub>2</sub>O (1) and [Dy(H<sub>2</sub>O)<sub>5</sub>(HMPA)<sub>2</sub>]I<sub>3</sub>·2HMPA (2) (HMPA = hexamethylphosphoramide), exhibit magnetic hysteresis which remains open up to 9 K for 1@Y and 10 K for 2@Y, respectively. *Ab initio* calculations, using step-by-step elimination of fragments to generate a series of model complexes, reveal that the secondary coordination sphere plays a key role in controlling the magnetisation reversal barrier and predict that the removal of outer-sphere molecules and anions will enhance the barrier further.**

Controlling the coordination environment/geometry at the level of a single metal ion is a challenging task for synthetic chemists, but of great importance for the miniaturisation of data storage.<sup>1</sup> Single-Molecule Magnets (SMMs) are magnetically independent molecules in which the magnetic moment resists reorientation upon removal of an external magnetising field, which results in the appearance of magnetic hysteresis of molecular origin.<sup>2</sup> These unique molecules are considered amongst the best candidates for advanced quantum-based devices (devices for data storage, quantum computing and spintronics).<sup>3</sup> Mononuclear SMMs (usually referred to as Single-Ion Magnets (SIMs)) comprise one or more ligands binding at a single metal site to control the coordination geometry.<sup>4</sup> The large single-ion magnetic anisotropy and unquenched orbital angular momentum of the f-elements have established them as enticing targets for the construction of SMMs/SIMs.<sup>5</sup> It has been recently demonstrated that for lanthanide-based SMMs, the prerequisites for large energy barriers for the reorientation of magnetisation ( $U_{\text{eff}}$ ) and magnetic hysteresis at higher temperatures are the

double degeneracy of the largest  $m_j$  ground substate and a large separation from the first and higher excited states.<sup>6</sup> According to computationally-demanding theoretical calculations (CASSCF-SO) and elegant electrostatic models, it has become clear that for improved dysprosium(III) based SMMs/SIMs, the type of coordination environment has to possess a strongly axial character.<sup>7</sup> In this way, shortcuts through the energy barrier can be suppressed and higher energy barriers are achieved for thermal relaxation.<sup>8</sup> Complexes with a high-order symmetry axis such as square antiprismatic ( $D_{4d}$ ),<sup>9</sup> trigonal bipyramidal ( $D_{3h}$ ),<sup>10</sup> sandwich type<sup>11</sup> (with cyclooctatetraene (COT) and cyclopentadienyl (Cp) anionic ligands) and pentagonal bipyramidal ( $D_{5h}$ ),<sup>12</sup> have been suggested as an effective way to favour slower relaxation by reducing transverse anisotropy, thus suppressing Quantum Tunnelling of the Magnetisation (QTM). Additionally higher-order symmetry also helps to align the  $g_{zz}$  axis of the excited state co-linear to the ground state, which leads to relaxation *via* higher excited states and hence larger magnetisation reversal barriers.<sup>13</sup>

Recently, we reported a trigonal bipyramidal Ni(II) complex with an unprecedented axial magnetic anisotropy by using bulky axial ligands.<sup>14</sup> Motivated by these results, we have investigated a wide range of bulky monodentate ligands to synthesise axial Dy(III) systems. Herein we report the synthesis, structure, magnetic properties and *ab initio* studies of [Dy(H<sub>2</sub>O)<sub>5</sub>(HMPA)<sub>2</sub>]Cl<sub>3</sub>·HMPA·H<sub>2</sub>O (1) and [Dy(H<sub>2</sub>O)<sub>5</sub>(HMPA)<sub>2</sub>]I<sub>3</sub>·2HMPA (2) {HMPA = hexamethylphosphoramide, [(CH<sub>3</sub>)<sub>2</sub>N]<sub>3</sub>PO}, which represent new examples of Dy(III) SMMs with close to ideal pentagonal bipyramidal geometry. We also report the diamagnetic yttrium analogues; [Y(H<sub>2</sub>O)<sub>5</sub>-(HMPA)<sub>2</sub>]Cl<sub>3</sub>·HMPA·H<sub>2</sub>O (3) and [Y(H<sub>2</sub>O)<sub>5</sub>(HMPA)<sub>2</sub>]I<sub>3</sub>·2HMPA (4) and the Dy-doped yttrium analogues (see ESI†).<sup>15</sup> We find that the strong axial crystal field promoted by the two axial HMPA ligands in complexes 1 and 2 leads to hysteresis loops that remain open until 9 K for the Dy-doped yttrium chloride analogue (1@Y), and 10 K for the Dy-doped yttrium iodide analogue (2@Y). In terms of a large barrier  $U_{\text{eff}}$  and magnetic hysteresis temperature, these complexes are amongst the best air-stable SMMs.

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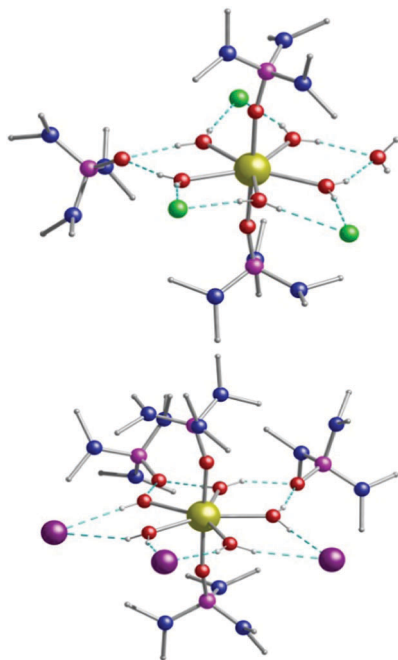


Fig. 1 The structure of **1** and **2**. Dy, gold; O, red; N, blue, P, pink; Cl, green; I, purple. Hydrogen atoms are omitted for clarity, except the H atoms of the equatorial water ligands and only one of the two crystallographically unique Dy units in **1** is shown.

The synthesis of **1** and **2** was accomplished by treatment of the hexamethylphosphoramide ligand (HMPA) with  $\text{DyCl}_3$  and  $\text{DyI}_3$ , respectively, in hot THF under aerobic conditions. Work-up and crystallization (see ESI<sup>†</sup>) from dichloromethane/toluene (1 : 1) produces large colourless prismatic crystals of  $[\text{Dy}(\text{H}_2\text{O})_5(\text{HMPA})_2]\text{Cl}_3 \cdot \text{HMPA} \cdot \text{H}_2\text{O}$  (**1**) and  $[\text{Dy}(\text{H}_2\text{O})_5(\text{HMPA})_2]\text{I}_3 \cdot 2\text{HMPA}$  (**2**) (Fig. 1). The phase purity of the bulk samples was confirmed by powder X-ray diffraction (Fig. S1–S4, ESI<sup>†</sup>).

The chloride analogue **1** crystallises in the monoclinic space group  $P2_1/c$  with the asymmetric unit containing two  $[\text{Dy}(\text{H}_2\text{O})_5(\text{HMPA})_2]\text{Cl}_3 \cdot \text{HMPA} \cdot \text{H}_2\text{O}$  units (Table S1 and Fig. S6, S7, ESI<sup>†</sup>). The seven coordinate  $\text{Dy}(\text{III})$  centre is in a pentagonal bipyramidal geometry ( $D_{5h}$  symmetry) for both crystallographically independent Dy centres, as confirmed *via* continuous shape measures analysis,<sup>16</sup> which provides an estimate of the distortion from the perfect polyhedron, giving a value of 0.154 for Dy1 (**1a**) and 0.284 for Dy2 (**1b**) (where 0 corresponds to the ideal structure) (Fig. S5 and Table S4, ESI<sup>†</sup>). Two HMPA ligands occupy the axial positions with five water molecules in the equatorial positions (Fig. S6, ESI<sup>†</sup>). The charge is balanced by three chloride anions in the second coordination sphere, with one HMPA and one  $\text{H}_2\text{O}$  co-crystallised (Fig. S8, ESI<sup>†</sup>). The axial Dy–O distances are 2.2171(14) and 2.2211(14) for Dy1 (**1a**) and 2.2096(15) and 2.2282(15) Å for Dy2 (**1b**) while the equatorial Dy–O distances fall in the range 2.3377(16)–2.3600(16) Å for Dy1 (**1a**) and 2.3365(16)–2.3587(17) Å for Dy2 (**1b**). The axial O–Dy–O angle is almost linear ( $\sim 176^\circ$ ) while the equatorial O–Dy–O angles range between  $71.41(6)^\circ$  and  $72.63(6)^\circ$ , close to the ideal angle of  $72^\circ$  (Table S2, ESI<sup>†</sup>). Extensive intermolecular hydrogen bonding is observed in **1** connecting all the moieties in the

structure (Fig. S8, ESI<sup>†</sup>). Complex  $[\text{Dy}(\text{H}_2\text{O})_5(\text{HMPA})_2]\text{I}_3 \cdot 2\text{HMPA}$  (**2**) is similar to **1** but crystallises in the monoclinic  $Cc$  space group. The asymmetric unit of **2** contains one  $\text{Dy}(\text{III})$  complex, three iodide anions and two co-crystallised HMPA ligands (Fig. S9, ESI<sup>†</sup>). We note that in complex **2**, the  $\text{Dy}(\text{III})$  pentagonal bipyramidal geometry is closer to the ideal polyhedron as confirmed *via* continuous shape measures analysis (0.131, Table S5, ESI<sup>†</sup>).<sup>16</sup> The axial Dy–O distances in **2** are 2.202(4) and 2.208(4) Å, while the equatorial Dy–O distances range between 2.343(4)–2.375(4) Å. The axial O–Dy–O angle is  $\sim 178^\circ$  with the equatorial O–Dy–O angles ranging from  $70.91(14)^\circ$  to  $73.18(14)^\circ$  (Table S3, ESI<sup>†</sup>). Additionally, analysis of the crystal packing reveals that between discrete molecules the nearest Dy···Dy distance for **1** is 9.026 Å, while for **2** the  $\text{Dy}(\text{III})$  centres are better isolated (11.767 Å) with only intramolecular hydrogen bonding observed (Fig. S10, ESI<sup>†</sup>).

The dc magnetic susceptibility measurements for complexes **1** and **2** are described in the ESI<sup>†</sup> (Fig. S11 and S12). The  $\chi_M T$  profiles (*i.e.* the practically unchanged  $\chi_M T$  product between 290–7 K and the sharp drop at low temperature) are similar to reported high performance Dy-SMM complexes.<sup>17</sup> The field cooled (FC) and zero-field cooled (ZFC) magnetic susceptibility (Fig. S13 and S14, ESI<sup>†</sup>) diverge at 7 K and 8 K for **1** and **2** respectively, while the maxima are observed at  $\sim 7$  K for both complexes, indicative of the magnetic blocking temperature,  $T_B$ .

The presence of magnetic blocking was further confirmed by field dependent magnetisation measurements,  $M(H)$  loops, carried out at different temperatures, with an average sweep rate of  $\sim 4 \text{ mT s}^{-1}$  (for a full cycle). Both **1** and **2** display clear magnetic hysteresis which remains open up to 6 K (9 K for **1**@Y) (Fig. S15–S17, ESI<sup>†</sup>) and 9 K (10 K for **2**@Y), respectively (Fig. 2 and Fig. S18, ESI<sup>†</sup>). The shape of the  $M(H)$  loops is strongly affected by the faster relaxation around zero field arising due to unsuppressed Quantum Tunnelling of the Magnetisation. The  $M(H)$  loops display a large coercivity of 1 T for **1**@Y and 1.1 T for **2**@Y at 2 K.

Ac susceptibility measurements between 1–1488 Hz in zero external dc field with a small oscillating ac field of  $H_{ac} = 3 \text{ Oe}$  were

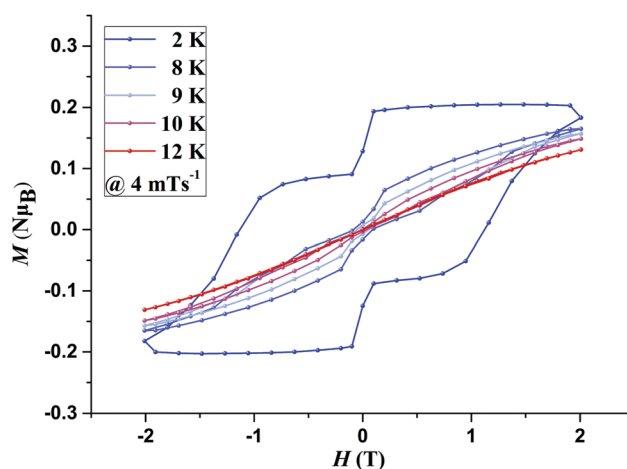


Fig. 2 Magnetic hysteresis measurements for **2**@Y with an average sweep rate of  $4 \text{ mT s}^{-1}$ .



performed to investigate the slow magnetic relaxation. The in phase  $\chi_M'$  and out of phase  $\chi_M''$  magnetic susceptibility data show frequency dependent peaks with well-defined maxima at 30 K for **1** (Fig. S19 and S20, ESI†) and 38 K for **2** (Fig. S21 and S22, ESI†) for a frequency of 1488 Hz, indicating high magnetisation reversal barriers. The plot of  $\chi_M''$  vs. frequency yielded temperature dependent maxima in the range of 17–30 K for **1** (Fig. S23, ESI†) and 17–37 K for **2** (Fig. S24, ESI†). The relaxation times,  $\tau$ , were extracted from the fittings of the Argand plots of  $\chi_M''$  vs.  $\chi_M'$  using the generalized Debye model.<sup>18</sup> The  $\alpha$  parameters found are in the range of 0.013–0.10 (17–30 K) for **1** (Fig. S25, ESI†) and 0.015–0.13 (17–37 K) for **2** (Fig. S26, ESI†), indicating a narrow distribution of relaxation times. The Arrhenius plots, constructed by using the extracted relaxation times,  $\tau$ , were fitted with the Arrhenius equation

$$\tau = \tau_0 \cdot \exp\left(\frac{U_{\text{eff}}}{T}\right).$$

The  $\ln \tau$  versus  $T^{-1}$  plots were fit considering only the Orbach process, to yield an energy barrier of  $U_{\text{eff}} = 460$  K,  $\tau_0 = 2 \times 10^{-11}$  s for **1** (Fig. S27, top, ESI†) and  $U_{\text{eff}} = 600$  K,  $\tau_0 = 1.2 \times 10^{-11}$  s for **2** (Fig. S27, bottom, ESI†). In order to further understand the mechanism that governs the magnetic relaxation observed for complexes **1** and **2**, we have performed *ab initio* calculations using the CASSCF/RASSI-SO/SINGLE\_ANISO approach implemented in MOLCAS8.0 (see ESI†). This methodology has a proven track record to rationalize, and also predict at times, a reliable mechanism of magnetic relaxation.<sup>12,13,19,20</sup> The eight Kramers Doublets (KDs), corresponding to the  ${}^6\text{H}_{15/2}$  state, span within energy ranges of  $882.1 \text{ cm}^{-1}$  and  $919.0 \text{ cm}^{-1}$ , respectively for complexes **1** and **2**. For both complexes,  $m_j = \pm 15/2$  is the ground state with  $m_j = \pm 13/2$  as the first excited state and  $m_j = \pm 1/2$  as the second-excited state. It is also worth mentioning here that the second excited state for complex **1** is mixed in nature whereas for complex **2**, it is pure  $m_j = \pm 1/2$  (Fig. 3 and Fig. S29, ESI†). The computed ground and first-excited state magnetic anisotropy for **1** and **2** are purely Ising in nature ( $g_{zz} \approx 19.97$  and  $17.10$  respectively, with  $g_{xx} = g_{yy} = 0$ , see Fig. S28 for the  $g_{zz}$  axis of **1** and **2**, ESI†) suggesting a very small operative QTM/TA-QTM within the ground and first-excited KDs (Fig. 3 and Fig. S29, Table S6, ESI†).

The Orbach process related to the ground and first excited state of opposite magnetisation is calculated to be very small (Fig. 3 and Fig. S29, ESI†). However, the second excited state for both complexes is found to have a substantial transverse magnetic component with a larger operative TA-QTM relaxation ( $\sim 0.5$  and  $3.5 \mu_B$  for complexes **1** and **2**, respectively). The computed  $g_{zz}$  axis for both the ground and first excited states is collinear (see Fig. S28, ESI†) in complexes **1** and **2** (less than  $2^\circ$ ). This suggests magnetic relaxation *via* the second excited KD ( $-15/2 \rightarrow -13/2 \rightarrow -1/2 \rightarrow +1/2 \rightarrow +13/2 \rightarrow +15/2$ ) with the  $U_{\text{cal}}$  value 584 K and 640 K for complexes **1** and **2**, respectively. In complex **2**, the calculated and the experimental energy barrier are in close agreement (640 K and 600 K respectively) whereas in complex **1**, the agreement is not as good (584 K and 460 K, respectively). This is expected as in complex **1**, the ground state QTM is larger compared to complex **2** (Fig. 3 and Fig. S29, ESI†). Furthermore, the structural parameters and LoProp<sup>22</sup> charge analysis (see Fig. S31, ESI†) suggest a stronger axial ligand field

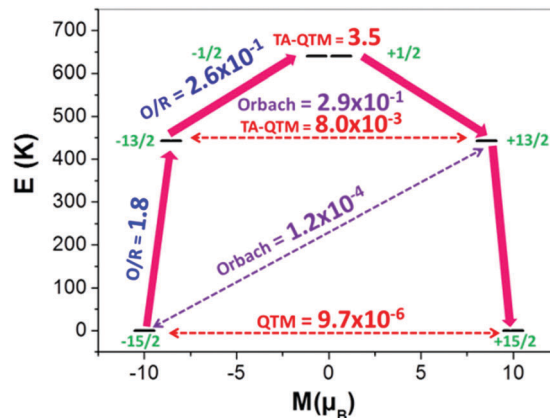


Fig. 3 *Ab initio* calculated relaxation dynamics for complex **2**. The arrows show the connected energy states with the number representing the matrix element of the transverse moment (see text for details). Here, QTM = quantum tunnelling of the magnetisation, TA-QTM = thermally assisted QTM, O/R = Orbach/Raman process. The numbers above each arrow represent corresponding transversal matrix elements for the transition magnetic moments as described in ref. 21 (see ESI† for further information).

in complex **2**. Although the LoProp charges on the pentagonal plane water molecules are larger in complex **2**, the even larger LoProp charges on  $\text{Cl}^-$  ions found in **1**, compared to  $\text{I}^-$  ions in **2**, result in a larger transverse ligand field in **1** (Fig. S31, ESI†). Notably, the non-axial term ( $B_k^q$ , where  $q \neq 0$  and  $k = 2, 4$ , and  $6$ ) to the axial term ( $B_k^0$ , where  $q = 0$  and  $k = 2, 4$ , and  $6$ ) ratio is found to be smaller for complex **2** compared to **1**, confirming the lower operational QTM and the larger  $U_{\text{cal}}$  value for **2** compared to **1** (Table S7, ESI†).<sup>20</sup> In an attempt to evaluate how strongly the  $U_{\text{cal}}$  values are related to molecules/fragments in the secondary coordination sphere, we have created a number of different model systems by step-by-step elimination of each fragment to generate a series of model complexes (Fig. 4 and Fig. S32, ESI†), with the last model possessing a mono-coordinated  $\text{Dy(III)}$  with one HMPA ligand. Importantly, these models highlight that the removal of anions/molecules in the secondary coordination sphere is able to enhance the computed value of  $U_{\text{cal}}$ . As these outer-sphere anions/molecules are removed, this causes a lowering of the charges in the equatorial region around the  $\text{Dy(III)}$  ion and thus increases the crystal field interaction along the axial direction (Fig. S31 and Table S8, ESI†). One can in fact attain a theoretical value of  $\sim 3100$  K for a two coordinate  $\{\text{Dy}(\text{HMPA})_2\}^{3+}$  model (Fig. 4 and Fig. S32, ESI†).<sup>7</sup> This can be attributed to the increase of the crystal field interaction along the axial direction as we move to lower-coordinate models (Fig. S34 and Table S8, ESI†). Furthermore, our calculations reveal that replacement of the halide anions, with large bulky anions, can cause a reduction in the charges in the equatorial region of the  $\text{Dy(III)}$  ion leading to larger barrier heights. Also, replacing small counter anions ( $\text{X}^-$ ) with large counter anions ( $\text{PF}_6^-$ ) could potentially decrease the crowding around the  $\text{Dy(III)}$  ion by eliminating secondary coordination sphere molecules or increasing their distance from the  $\text{Dy(III)}$  ion and hence result in a lowering of the charges in the equatorial plane. Collectively both effects are expected to enhance significantly the barrier height and efforts are currently underway to achieve this in our laboratory.



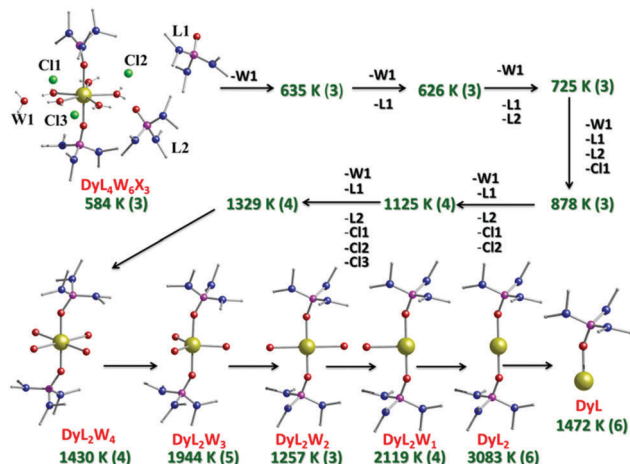


Fig. 4 *Ab initio* SINGLE\_ANISO computed  $U_{\text{cal}}$  values (in K, green text) for all the models created by stepwise removal of fragments for complex **1**. The bold number in brackets for each model gives the Kramers doublet through which relaxation is expected (ground state = KD (1); first excited state = KD (2) etc.). For more information see Fig. S32 and S33 (ESI<sup>†</sup>).

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## Conflicts of interest

There are no conflicts to declare.

## Notes and references

- R. Vincent, S. Klyatskaya, M. Ruben, W. Wernsdorfer and F. Balestro, *Nature*, 2012, **488**, 357.
- J. Lu, M. Guo and J. Tang, *Chem. – Asian J.*, 2017, **12**, 2772; S. G. McAdams, A.-M. Ariciu, A. K. Kostopoulos, J. P. S. Walsh and F. Tuna, *Coord. Chem. Rev.*, 2017, **346**, 216; F. Pointillart, O. Cadot, B. Le Guennic and L. Ouahab, *Coord. Chem. Rev.*, 2017, **346**, 150; S. Demir, I.-R. Jeon, J. R. Long and T. D. Harris, *Coord. Chem. Rev.*, 2015, **289**, 149; Y.-S. Meng, S.-D. Jiang, B.-W. Wang and S. Gao, *Acc. Chem. Res.*, 2016, **49**, 2381; D. N. Woodruff, R. E. P. Winpenny and R. A. Layfield, *Chem. Rev.*, 2013, **113**, 5110.
- G. Aromi, D. Aguila, P. Gamez, F. Luisic and O. Roubeau, *Chem. Soc. Rev.*, 2012, **41**, 537; C. Godfrin, A. Ferhat, R. Ballou, S. Klyatskaya, M. Ruben, W. Wernsdorfer and F. Balestro, *Phys. Rev. Lett.*, 2017, **119**, 187702; S. Thiele, F. Balestro, R. Ballou, S. Klyatskaya, M. Ruben and W. Wernsdorfer, *Science*, 2014, **344**, 1135; L. Bogani and W. Wernsdorfer, *Nat. Mater.*, 2008, **7**, 179; M. Affronte, *J. Mater. Chem.*, 2009, **19**, 1731.
- S. Gomez-Coca, E. Cremades, N. Aliaga-Alcald and E. Ruiz, *J. Am. Chem. Soc.*, 2013, **135**, 7010.
- C. Benelli and D. Gatteschi, *Introduction to Molecular Magnetism. From Transition Metals to Lanthanides*, Wiley-VCH, Weinheim, 2015;
- R. A. Layfield and M. Murugesu, *Lanthanides and Actinides in Molecular Magnetism*, Wiley-VCH, Weinheim, 2015.
- S. T. Liddle and J. van Slageren, *Chem. Soc. Rev.*, 2015, **44**, 6655; M. Gregson, N. F. Chilton, A.-M. Ariciu, F. Tuna, I. F. Crowe, W. Lewis, A. J. Blake, D. Collison, E. J. L. McInnes, R. E. P. Winpenny and S. T. Liddle, *Chem. Sci.*, 2016, **7**, 155.
- J. D. Rinehart and J. R. Long, *Chem. Sci.*, 2011, **2**, 2078; L. Ungur and L. F. Chibotaru, *Phys. Chem. Chem. Phys.*, 2011, **13**, 20086; S. K. Singh, T. Gupta and G. Rajaraman, *Inorg. Chem.*, 2014, **53**, 10835; L. Ungur and L. F. Chibotaru, *Inorg. Chem.*, 2016, **55**, 10043.
- T. Pugh, N. F. Chilton and R. A. Layfield, *Angew. Chem., Int. Ed.*, 2016, **55**, 11082; W.-B. Sun, P.-F. Yan, S.-D. Jiang, B.-W. Wang, Y.-Q. Zhang, H.-F. Li, P. Chen, Z.-M. Wang and S. Gao, *Chem. Sci.*, 2016, **7**, 684; P. Zhang, M. Perfetti, M. Kern, P. P. Hallmen, L. Ungur, S. Lenz, M. R. Ringenberg, W. Frey, H. Stoll, G. Rauhut and J. van Slageren, *Chem. Sci.*, 2018, **9**, 1221; L. Norel, L. E. Darago, B. Le Guennic, K. Chakarawet, M. I. Gonzalez, J. H. Olshansky, S. Rigaut and J. R. Long, *Angew. Chem., Int. Ed.*, 2018, **57**, 1; S. Demir, M. Gonzalez, L. E. Darago, W. J. Evans and J. R. Long, *Nat. Commun.*, 2017, **8**, 2144; B. S. Dolinar, D. I. Alexandropoulos, K. R. Vignesh, T. A. James and K. R. Dunbar, *J. Am. Chem. Soc.*, 2018, **140**, 908.
- L. Sorace, C. Benelli and D. Gatteschi, *Chem. Soc. Rev.*, 2011, **40**, 3092; P.-E. Car, M. Perfetti, M. Mannini, A. Favre, A. Caneschi and R. Sessoli, *Chem. Commun.*, 2011, **47**, 3751; M. A. Aldamen, J. M. Clemente-Juan, E. Coronado, C. Marti-Gastaldo and A. Gaita-Arriño, *J. Am. Chem. Soc.*, 2008, **130**, 8874.
- P. Zhang, L. Zhang, C. Wang, S. Xue, S.-Y. Lin and J. Tang, *J. Am. Chem. Soc.*, 2014, **136**, 4484; K. L. M. Harriman, J. L. Brosmer, L. Ungur, P. L. Diaconescu and M. Murugesu, *J. Am. Chem. Soc.*, 2017, **139**, 1420.
- S.-D. Jiang, B.-W. Wang, H.-L. Sun, Z.-M. Wang and S. Gao, *J. Am. Chem. Soc.*, 2011, **133**, 4730; F.-S. Guo, B. M. Day, Y.-C. Chen, M.-L. Tong, A. Mansikkamäki and R. A. Layfield, *Angew. Chem., Int. Ed.*, 2017, **56**, 11445; C. A. P. Goodwin, F. Ortu, D. Reta, N. F. Chilton and D. P. Mills, *Nature*, 2017, **548**, 439.
- Y.-C. Chen, J.-L. Liu, L. Ungur, J. Liu, Q.-W. Li, L.-F. Wang, Z.-P. Ni, L. F. Chibotaru, X.-M. Chen and M.-L. Tong, *J. Am. Chem. Soc.*, 2016, **138**, 2829; Y.-S. Ding, N. F. Chilton, R. E. P. Winpenny and Y.-Z. Zheng, *Angew. Chem., Int. Ed.*, 2016, **55**, 16071; J. Liu, Y.-C. Chen, J.-L. Liu, V. Vieru, L. Ungur, J.-H. Jia, L. F. Chibotaru, Y. Lan, W. Wernsdorfer, S. Gao, X.-M. Chen and M.-L. Tong, *J. Am. Chem. Soc.*, 2016, **138**, 5441; S. K. Gupta, T. Rajeshkumar, G. Rajaraman and R. Murugavel, *Chem. Sci.*, 2016, **7**, 5181–5191; J. L. Liu, Y. C. Chen, Y. Z. Zheng, W. Q. Lin, L. Ungur, W. Wernsdorfer, L. F. Chibotaru and M. L. Tong, *Chem. Sci.*, 2013, **4**, 3310.
- S. K. Singh, T. Gupta, M. Shanmugam and G. Rajaraman, *Chem. Commun.*, 2014, **50**, 15513.
- K. E. R. Marriott, L. Bhaskaran, C. Wilson, M. Medarde, S. T. Ochsenbein, S. Hill and M. Murrie, *Chem. Sci.*, 2015, **6**, 6823.
- Z. Hou, K. Kobayashi and H. Yamazaki, *Chem. Lett.*, 1991, 265; T. Imamoto, Y. Yamanoi, H. Tsuruta, K. Yamaguchi, M. Yamazaki and J. Inanaga, *Chem. Lett.*, 1995, 949.
- M. Llune, D. Casanova, J. Cirera, J. Bo, P. Alemany and S. Alvarez, *SHAPE (Version 2.1)*, Barcelona, 2013.
- Z. Zhu, M. Guo, X.-L. Li and J. Tang, *Coord. Chem. Rev.*, 2017, DOI: 10.1016/j.ccr.2017.10.030.
- D. Gatteschi, R. Sessoli and J. Villain, *Molecular Nanomagnets*, Oxford Univ. Press, 2006.
- M. K. Singh, N. Yadav and G. Rajaraman, *Chem. Commun.*, 2015, **51**, 17732; M. K. Singh and G. Rajaraman, *Chem. Commun.*, 2016, **52**, 14047.
- A. K. Bar, P. Kalita, M. K. Singh, G. Rajaraman and V. Chandrasekhar, *Coord. Chem. Rev.*, 2018, **367**, 163.
- L. Ungur, M. Thewissen, J. P. Costes, W. Wernsdorfer and L. F. Chibotaru, *Inorg. Chem.*, 2013, **52**, 6328.
- L. Gagliardi, R. Lindh and G. Karlström, *J. Chem. Phys.*, 2004, **121**, 4494.

